# Measurements and Correlation of High-Pressure Densities of Imidazolium-Based Ionic Liquids

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In the present work, experimental density measurements are reported along with the derived thermodynamic properties, such as the isothermal compressibility ( $\kappa_T$ ), the isobaric expansivity ( $\alpha_p$ ), and the thermal pressure coefficient ( $\gamma_v$ ) for imidazolium-based ionic liquids (ILs), namely, 1-ethyl-3-methylimidazolium methylsulfate [C<sub>2</sub>mim][MeSO<sub>4</sub>], 1-ethyl-3-methylimidazolium ethylsulfate [C<sub>2</sub>mim][EtSO<sub>4</sub>], 1,3-diethylimidazolium bis-(trifluoromethylsulfonyl)imide [C<sub>2</sub>eim][Tf<sub>2</sub>N], and 1-decyl-3-methylimidazolium bis(trifluoromethylsulfonyl)imide [C<sub>10</sub>mim][Tf<sub>2</sub>N] in the pressure (0.10 < *P*/MPa < 35.00) and temperature (293.15 < *T*/K < 393.15) domains. It is shown that experimental densities are in good agreement with the predicted densities obtained by the Gardas and Coutinho method and the correlation using the Tait equation and Sanchez–Lacombe equation of state.

# Introduction

Ionic liquids (ILs), in general, are the combination of large organic cations with various alkyl substituents and inorganic or organic anions, so that a very large number of ILs are possible. To better understand the nature of ionic liquids and rationally expand their applications, knowledge of their physical properties is required. The design of industrial processes and new products based on ILs can only be achieved when their thermophysical properties are adequately characterized.

High-pressure density data of pure ionic liquids available in the literature<sup>1-22</sup> are summarized in Table 1. This paper is a continuation of previous works<sup>8,10,18</sup> on high-pressure densities and derived thermodynamic properties of ILs. In this work, the experimental measurements of the pressure ( $0.10 \le P/MPa \le$ 35.00) and temperature (293.15 < T/K < 393.15) dependence of the density of imidazolium-based ionic liquids (ILs), namely, 1-ethyl-3-methylimidazolium methylsulfate [C<sub>2</sub>mim][MeSO<sub>4</sub>], 1-ethyl-3-methylimidazolium ethylsulfate [C<sub>2</sub>mim][EtSO<sub>4</sub>], 1,3diethylimidazolium bis(trifluoromethylsulfonyl)imide [C2eim]-[Tf<sub>2</sub>N], and 1-decyl-3-methylimidazolium bis(trifluoromethylsulfonyl)imide [C<sub>10</sub>mim][Tf<sub>2</sub>N], are presented. The experimental densities were correlated with the Tait equation<sup>23</sup> and the Sanchez-Lacombe equation of state (S-L EoS)<sup>24</sup> and compared against the estimation method proposed by Gardas and Coutinho.<sup>25</sup> Derived thermodynamic properties, such as the isothermal compressibility ( $\kappa_T$ ), the isobaric expansivity ( $\alpha_p$ ), and the thermal pressure coefficient  $(\gamma_n)$  are reported as Supporting Information.

The objective of this work is to contribute to the databank of thermodynamic properties of pure ILs and to investigate the relationship between densities and the ionic structures, in order to establish principles for the molecular design of ILs. For that purpose, the [ $C_2$ mim] cation was studied in combination with

two anions,  $[MeSO_4]$  and  $[EtSO_4]$ , to conclude about the anion effect. The volumetric study of  $[C_2eim][Tf_2N]$  and  $[C_{10}mim]$ - $[Tf_2N]$  is made possible by the comparison with our previous study,<sup>10</sup> to understand the effect of alkyl chain length on the imidazolium ring on the density and derived properties. The effect of alkyl chain length on the imidazolium-based ILs, having the  $[Tf_2N]$  anion, on the density and derived properties has been studied.

# **Experimental Section**

*Materials.* Experimental densities were measured for imidazolium-based ILs, namely,  $[C_2mim][MeSO_4]$ ,  $[C_2mim][EtSO_4]$ ,  $[C_2eim][Tf_2N]$ , and  $[C_{10}mim][Tf_2N]$ . All ILs were acquired from IoLiTec (Ionic Liquids Technologies GmbH & Co. KG, Germany) with mass fraction purities > 99 % and a bromide impurity mass fraction <  $10^{-4}$ .

To reduce the water content and volatile compounds to negligible values, a vacuum (0.1 Pa) at moderate temperature (353 K) for at least 48 h was applied to all the IL samples prior to their use. The purity of each IL was checked by <sup>1</sup>H, <sup>13</sup>C, and <sup>19</sup>F NMR, and the water content in the ILs was determined with a Metrohm 831 Karl Fischer coulometer indicating very low levels of water mass fraction content, as (136, 191, 95, and 118)  $\cdot 10^{-6}$  for [C<sub>2</sub>mim][MeSO<sub>4</sub>], [C<sub>2</sub>mim][EtSO<sub>4</sub>], [C<sub>2</sub>eim]-[Tf<sub>2</sub>N], and [C<sub>10</sub>mim][Tf<sub>2</sub>N], respectively. The analyte used for the coulometric Karl Fischer titration was Hydranal-Coulomat AG from Riedel-de Haën.

**Experimental Procedure.** Experimental densities were measured using an Anton Paar DMA 60 digital vibrating tube densimeter, with a DMA 512P measuring cell in the temperature range (293.15 to 393.15) K and pressure range (0.10 to 35.00) MPa. The temperature in the vibrating tube cell was measured with a platinum resistance probe which has a temperature uncertainty of  $\pm$  0.01 K coupled with a GW instek dual display digital multimeter GDM-845. A Julabo P-5 thermostatic bath with silicone oil as circulating fluid was used in

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		maximum	data			
ionic liquid	T/K	<i>P/MPa</i>	method <sup>a</sup>	points	ref	
	298.2 to 343.2	200	U	20	1	
$[C_4 mim][PF_6]$	298.2 to 343.2 298.15 to 323.15	200 100	U V	20 144	2	
	298.15 to 323.15 293.15 to 353.15	20	v D	20	3	
	298.15 to 398.15	40	V	45	4	
	293.48 to 414.94	40	v	46	5	
	294.9 to 335.1	20	D	09	6	
	312.8 to 472.3	200	D	189	7	
[C <sub>6</sub> mim][PF <sub>6</sub> ]	293.15 to 393.15	10	V	77	8	
[060000][116]	294.9 to 335.2	20	D	09	6	
	293.15 to 353.15	20	D	20	9	
[C <sub>8</sub> mim][PF <sub>6</sub> ]	298.2 to 343.2	202	U	20	1	
[08][11.6]	293.15 to 393.15	10	v	77	8	
	295.1 to 335.2	20	D	09	6	
	293.15 to 353.15	20	D	20	9	
[C <sub>2</sub> mim][BF <sub>4</sub> ]	293.15 to 393.15	30	V	96	10	
$[C_4 \text{mim}][BF_4]$	278.15 to 333.15	60	V	26	11	
[ 4	298.34 to 332.73	60	V	67	2	
	293.15 to 353.15	20	D	20	3	
	298.15 to 398.15	40	V	45	4	
	293.15 to 393.15	10	V	77	8	
	293.49 to 414.93	40	V	46	5	
	273.15 to 348.15	300	V	62	12	
	313.1 to 472.2	200	D	189	7	
	298.15 to 398.15	40	V	45	13	
[C <sub>8</sub> mim][BF <sub>4</sub> ]	298.2 to 343.2	207	U	21	1	
	293.15 to 393.15	10	V	77	8	
[C <sub>4</sub> mmim][BF <sub>4</sub> ]	313.15 to 393.15	10	V	63	8	
[C <sub>2</sub> mim][Tf <sub>2</sub> N]	293.49 to 414.93	40	V	46	5	
	293.15 to 393.15	30	V	96	10	
[C <sub>3</sub> mim][Tf <sub>2</sub> N]	298.15 to 333.15	60	V	165	14	
[C <sub>4</sub> mim][Tf <sub>2</sub> N]	298.15 to 328.20	60	V	168	15	
	293.49 to 414.92	40	V	46	5	
[C <sub>5</sub> mim][Tf <sub>2</sub> N]	298.15 to 333.15	60	V	165	14	
[C <sub>6</sub> mim][Tf <sub>2</sub> N]	298.15 to 328.20	60	V	156	15	
	298.15 to 423.15	40	V	28	16	
	293.15 to 338.15	65	V	160	17	
[C <sub>7</sub> mim][Tf <sub>2</sub> N]	293.15 to 393.15	30	V	96	10	
[C <sub>8</sub> mim][Tf <sub>2</sub> N]	293.15 to 393.15	30	V	96	10	
[C2mim][CF3SO3]	293.15 to 393.15	35	V	91	18	
$[C_4 mim][CF_3SO_3]$	293.15 to 393.15	10	V	77	8	
$[C_4 mim][C(CN)_3]$	293.15 to 393.15	30	V	96	12	
[C <sub>1</sub> mim][MeSO <sub>4</sub> ]	313.15 to 333.15	25	V	34	19	
[C <sub>2</sub> mim][EtSO <sub>4</sub> ]	292.49 to 414.95	40	V	46	5	
10 1 10 00 1	283.15 to 333.15	35	V	63	20	
[C <sub>4</sub> mim][OcSO <sub>4</sub> ]	298.15 to 428.15	60	V	214	21	
	312.9 to 472.6	200	D	178	7	
[THTDP][Tf <sub>2</sub> N]	298.15 to 333.43	65	V	126	22	
[THTDP][Cl]	298.13 to 333.14	65	V	134	22	
[THTDP][CH <sub>3</sub> COO]	298.13 to 334.11	65	V	144	22	
$[C_4 py][BF_4]$	298.2 to 343.2	204	U	20	1	
[C <sub>3</sub> mpy][Tf <sub>2</sub> N]	293.15 to 393.15	35	V	91	18	
$[C_3 mpyr][Tf_2N]$	293.15 to 393.15	35	V	91 01	18	
[C <sub>4</sub> mpyr][Tf <sub>2</sub> N]	293.15 to 393.15	35	V V	91 01	18	
[C <sub>3</sub> mpip][Tf <sub>2</sub> N]	293.15 to 393.15	35	v	91	18	

Table 1. High-Pressure Density Data of Pure Ionic Liquids **Reported in the Literature** 

in a previous work of Gardas et al.<sup>26</sup> The equation proposed by Niesen<sup>27</sup> which has a solid theoretical basis as discussed by Holcom and Outcalt<sup>28</sup> was used for that purpose. The standard deviation of the fitting is in the order of  $\pm 1 \text{ kg} \cdot \text{m}^{-3}$ . The average absolute relative deviation, AARD, from the experimental data to the fitting is 0.06 % for all the ILs studied.

The calibration of the vibrating tube densimeter was described

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The influence of the viscosity on the densities was evaluated. To check the effect of viscosity in the density, a viscosity correction for compounds with viscosities < 100 mPa·s was applied with the equation proposed for the density uncertainty of an Anton Paar DMA 512 densimeter.<sup>29</sup> For compounds with viscosities higher than 400 mPa·s, the correction factor becomes constant<sup>30</sup> and equal to 0.5 kg $\cdot$ m<sup>-3</sup>, and between (100 and 400) mPa·s, the viscosity correction follows an intermediate behavior. Considering, for example, the available viscosity data for  $[C_2 mim][EtSO_4]^{31-33}$  at atmospheric pressure and temperature in the range (298.15 to 343.15) K, where the viscosity of ILs is < 100 mPa·s, allowed us to calculate an average density uncertainty less than 0.5 kg·m<sup>-3</sup>. For other ILs and/or other higher pressures where the viscosity increases, the correction value was assumed as 0.5 kg·m<sup>-3</sup>, being less than the uncertainty in the overall density data, which is 1 kg $\cdot$ m<sup>-3</sup>, and for that reason viscosity corrections were neglected in the present work.

# **Results and Discussion**

Density measurements were carried out at temperatures ranging from (293.15 to 393.15) K and pressures from (0.10 to 35.00) MPa. The experimental pressure-volume-temperature (PVT) data obtained are reported in Table 2 for all the ILs studied.

Density data for [C<sub>2</sub>mim][EtSO<sub>4</sub>] are already available in the open literature but at atmospheric pressure  $^{31-35}$  except the data by Jacquemin et al.<sup>5</sup> and Hofman et al.<sup>20</sup> The relative deviations between the experimental and literature density data of  $[C_2 mim][EtSO_4]$  are presented in Figure 1. Experimental data for [C<sub>2</sub>mim][EtSO<sub>4</sub>] are in good agreement with available literature values, the relative deviations ranging from (-0.64)to 1.27) % at atmospheric pressure and up to 2 % at high pressures. It is worth noting that the experimental density data of  $[C_2 mim][EtSO_4]$  are in excellent agreement (within 0.15 %) with the values of Krummen et al.,<sup>34</sup> while they are higher (within 1.27 %) than the values of Gomez et al.,<sup>30</sup> Rodriguez et al.,<sup>32</sup> and Gonzalez et al.<sup>33</sup> The relative deviations from the values of Jacquemin et al.<sup>5</sup> show both negative and positive values (ranging from -0.64 % to 0.48 %) at atmospheric pressure and (ranging from -2.00 % to 0.36 %) at higher pressure. The literature values of densities for this analysis were obtained using the Tait equation proposed by the authors.<sup>5</sup> The experimental density data of [C<sub>2</sub>mim][EtSO<sub>4</sub>] are in good agreement (percentage deviations ranging from 0.47 % to 0.71 %) with the values of Hofman et al.<sup>20</sup> These deviations can be due essentially to the purity of IL including water and halide content and also from the experimental technique adopted. To the best of our knowledge, no literature data on densities were previously available for [C<sub>2</sub>mim][MeSO<sub>4</sub>], [C<sub>2</sub>eim][Tf<sub>2</sub>N], and  $[C_{10}mim][Tf_2N].$ 

From the experimental densities for ILs having the  $[Tf_2N]^$ anion (comparison of [C<sub>2</sub>eim][Tf<sub>2</sub>N] and [C<sub>10</sub>mim][Tf<sub>2</sub>N]) and by the comparison of densities of [C<sub>2</sub>eim][Tf<sub>2</sub>N] with that of [C<sub>2</sub>mim][Tf<sub>2</sub>N] from our previous paper,<sup>10</sup> it is observed that as the alkyl chain length in the imidazolium cation increases the density of the corresponding IL decreases. Figure 2 shows

<sup>a</sup> D, dilatometer; V, vibrating tube density meter; U, ultrahighpressure density apparatus.

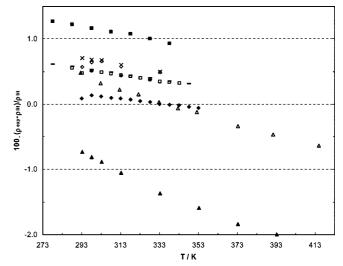
the thermostat circuit of the measuring cell which was held constant to  $\pm 0.01$  K.

The required pressure was generated and controlled with a Pressure Generator model 50-6-15, High Pressure Equipment Co., using acetone as the hydraulic fluid. The diameter of metallic tube is  $1.59 \cdot 10^{-3}$  m, and the buffer is more than 1 m long which guarantees the inexistence of diffusion of the hydraulic liquid in the liquid contained in the cell of the densimeter. Pressures were measured with a pressure transducer (Wika Transmitter S-10, WIKA Alexander Wiegand GmbH & Co.) with a maximum uncertainty of  $\pm$  0.03 MPa. An NI PCI-6220 data acquisition board (DAQ) from National Instruments (NI) was used for the real time registration of values of period, temperature, and pressure. Modules of temperature (NI SCC-FT01) and pressure (NI SCC-CI20) were installed into a NI SC-2345 carrier and connected to the DAQ board.

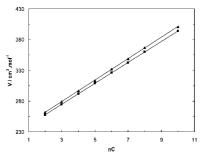
Table 2. Experimental Density,  $\rho$ , Data for [C<sub>2</sub>mim][MeSO<sub>4</sub>], [C<sub>2</sub>mim][EtSO<sub>4</sub>], [C<sub>2</sub>eim][Tf<sub>2</sub>N], and [C<sub>10</sub>mim][Tf<sub>2</sub>N] as a Function of Temperature and Pressure

		$ ho/kg \cdot m^{-3}$ at $T/K$							
<i>p</i> /MPa	293.15	298.15	303.15	313.15	333.15	353.15	373.15	393.15	
[C <sub>2</sub> mim][MeSO <sub>4</sub> ]									
0.10	1292.1	1288.0	1283.3	1276.6	1260.9	1245.5	1231.4	1218.1	
2.50	1293.0	1289.1	1284.3	1277.8	1262.3	1247.5	1232.8	1219.2	
5.00	1294.0	1290.1	1285.4	1278.9	1263.4	1248.7	1234.0	1220.5	
7.50	1295.0	1291.1	1286.4	1279.9	1264.5	1249.9	1235.2	1221.8	
10.00	1296.0	1292.1	1287.4	1281.0	1265.6	1251.1	1236.4	1223.1	
15.00	1298.0	1294.1	1289.5	1283.0	1267.8	1253.4	1238.8	1225.5	
20.00	1299.9	1296.1	1291.4	1285.1	1269.9	1255.6	1241.2	1227.9	
25.00	1301.8	1298.0	1293.4	1287.1	1272.0	1257.8	1243.5	1230.3	
30.00	1303.7	1300.0	1295.3	1289.1	1274.0	1260.0	1245.7	1232.6	
35.00	1305.6	1301.9	1297.3	1291.1	1276.1	1262.2	1248.0	1234.9	
0.10	1016.0		- 2	mim][EtS		1001-	1100 -		
0.10	1246.9	1244.3	1240.8	1233.1	1218.3	1204.2	1190.1	1176.6	
2.50	1249.0	1245.2	1241.6	1234.0	1219.2	1205.3	1191.3	1178.0	
5.00	1250.2	1246.3	1242.7	1235.1	1220.4	1206.6	1192.6	1179.3	
7.50	1251.3 1252.3	1247.4 1248.4	1243.8 1244.9	1236.2 1237.3	1221.6	1207.8	1193.9 1195.2	1180.7	
10.00	1252.5	1248.4 1250.6	1244.9	1237.5	1222.7 1225.0	1209.0 1211.4	1195.2	1182.0	
15.00 20.00	1254.5	1250.6	1247.0	1239.5	1225.0	1211.4 1213.8	1200.2	1184.6 1187.2	
20.00	1256.5	1252.7	1249.1	1241.0	1227.2	1215.8	1200.2	1187.2	
30.00	1258.5	1254.8	1251.1	1245.7	1229.5	1210.1	1202.0	1189.7	
35.00	1260.3	1258.7	1255.2	1245.8	1231.5	1218.4	1204.9	1192.2	
[C <sub>2</sub> eim][Tf <sub>2</sub> N]									
0.10	1469.2	1463.2	1457.7	1447.3	1425.7	1404.5	1384.1	1364.2	
2.50	1471.1	1465.5	1460.1	1449.6	1428.2	1407.5	1386.9	1367.2	
5.00	1473.1	1467.6	1462.1	1451.6	1430.4	1409.9	1389.5	1369.9	
7.50	1475.0	1469.5	1463.9	1453.7	1432.6	1412.2	1391.9	1372.5	
10.00	1476.9	1471.4	1466.1	1455.8	1434.8	1414.5	1394.4	1375.1	
15.00	1480.8	1475.3	1469.8	1459.7	1439.1	1419.0	1399.1	1380.2	
20.00	1484.2	1479.0	1473.6	1463.5	1443.2	1423.4	1403.7	1384.9	
25.00	1487.9	1482.5	1477.1	1467.2	1447.2	1427.5	1408.1	1389.6	
30.00	1491.5	1486.1	1480.9	1470.9	1450.9	1431.6	1412.4	1394.3	
35.00	1495.2	1489.5	1484.3	1474.6	1454.8	1435.6	1416.7	1398.6	
			[C <sub>1</sub>	omim][Ti	$f_2N$ ]				
0.10	1282.4	1278.0	1273.6	1265.3	1245.5	1228.6	1209.6	1191.8	
2.50	1285.6	1280.8	1276.4	1267.3	1247.7	1230.9	1212.4	1195.8	
5.00	1287.4	1282.7	1278.3	1269.3	1249.8	1233.2	1214.8	1198.3	
7.50	1289.3	1284.6	1280.3	1271.3	1251.9	1235.5	1217.3	1200.9	
10.00	1291.1	1286.4	1282.1	1273.2	1254.0	1237.6	1219.6	1203.4	
15.00	1294.6	1290.0	1285.8	1277.0	1258.0	1242.0	1224.2	1208.2	
20.00	1298.1	1293.5	1289.4	1280.7	1261.9	1246.2	1228.6	1212.8	
25.00	1301.4	1297.0	1292.9	1284.2	1265.7	1250.2	1232.9	1217.3	
30.00	1304.7	1300.3	1296.3	1287.7	1269.4	1254.1	1237.0	1221.7	
35.00	1308.0	1303.6	1299.6	1291.1	1273.0	1257.9	1241.1	1225.9	

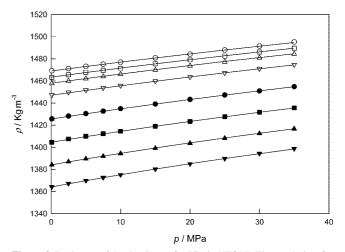
a comparison of the molar volumes obtained for the series of [C<sub>n</sub>mim][Tf<sub>2</sub>N] (this work and refs 10, 14, and 15) at two temperatures, (298.15 and 323.15) K, indicating linear trends with slopes in close agreement with the one anticipated by Esperança et al.<sup>14</sup> and Gardas et al.<sup>10</sup> The variation of the molar volume per addition of one carbon atom in the alkyl chain  $(\partial V_m/$  $\partial n$ ) is, respectively, [(17.1 ± 0.5) and (17.4 ± 0.5)] cm<sup>3</sup>·mol<sup>-1</sup> at temperatures (298.15 and 323.15) K. The change in molar volume by the addition of the  $-CH_2$  group, in  $[C_2mim][Tf_2N]$ to give  $[C_2 eim][Tf_2N]$ , is obtained as 19.6 cm<sup>3</sup>·mol<sup>-1</sup> at 298.15 K which is slightly higher (12.8 %) than that observed for a series of [C<sub>n</sub>mim][Tf<sub>2</sub>N] ILs. Inversely, as expected, for ILs having a similar cation, e.g.,  $[C_2 mim]^+$ , the change in molar volume by the addition of a -CH<sub>2</sub> group in the anion of IL  $[C_2 \text{mim}][MeSO_4]$  to give  $[C_2 \text{mim}][EtSO_4]$  is 17.3 cm<sup>3</sup>·mol<sup>-1</sup> at 298.15 K which is similar to that observed for the addition of a  $-CH_2$  group in the cation for a series of  $[C_n mim][Tf_2N]$ ILs. The molar volumes of  $[C_2 mim]^+$  increase with the effective anion size from [MeSO<sub>4</sub>] < [EtSO<sub>4</sub>] and by taking into account the results of our previous papers<sup>10</sup> show the trend  $[BF_4] <$  $[MeSO_4] < [EtSO_4] < [Tf_2N].$ 



**Figure 1.** Relative deviations between the experimental density data of  $[C_2mim][EtSO_4]$  and those reported in the literature as a function of temperature at 0.10 MPa:  $\Delta$ , Jacquemin et al.;<sup>5</sup>  $\diamond$ , Hofman et al.;<sup>20</sup>  $\Box$ , Gomez et al.;<sup>31</sup> –, Rodriguez et al.;<sup>32</sup> +, Gonzalez et al.;<sup>33</sup>  $\blacklozenge$ , Krummen et al.;<sup>34</sup> , Yang et al.<sup>35</sup> At 35.00 MPa:  $\blacktriangle$ , Jacquemin et al.;<sup>5</sup> ×, Hofman et al.<sup>20</sup>



**Figure 2.** Molar volumes at P = 0.1 MPa as a function of the number of carbons, *n*C, on the series of  $[C_nmim][NTf_2]$  (this work and references 10, 14, and 15): **I**, 298.15 K; **A**, 323.15 K.



**Figure 3.** Isotherms of the density,  $\rho$ , for  $[C_2 \text{cim}][\text{Tf}_2\text{N}]$ . The symbols refer to the experimental data:  $\bigcirc$ , 293.15 K;  $\square$ , 298.15 K;  $\triangle$ , 303.15 K;  $\bigtriangledown$ , 313.15 K;  $\blacklozenge$ , 333.15 K;  $\blacksquare$ , 353.15 K;  $\blacktriangle$ , 373.15 K;  $\blacktriangledown$ , 393.15 K. The curves are calculated with the Tait equation, eq 1.

### Modeling

**Tait Equation Correlation.** The liquid densities were correlated with the Tait equation,<sup>23</sup> and other thermodynamic properties such as the isothermal compressibility,  $\kappa_T$ , the isobaric expansivity,  $\alpha_p$ , and the thermal pressure coefficient,  $\gamma_v$ , were calculated and reported as Supporting Information. The results

Table 3. Coefficients of Equation 2, along with Standard Deviation of the Fit  $(\sigma)$ 

	$a_1$	$a_2$	$a_3 \cdot 10^4$	σ
ionic liquid	$kg \cdot m^{-3}$	$\overline{\mathrm{kg}}\cdot\mathrm{m}^{-3}\cdot\mathrm{K}^{-1}$	$\overline{\mathrm{kg}}\cdot\mathrm{m}^{-3}\cdot\mathrm{K}^{-2}$	$kg \cdot m^{-3}$
[C <sub>2</sub> mim][MeSO <sub>4</sub> ]	1593.502	-1.2430	7.3195	0.42
[C <sub>2</sub> mim][EtSO <sub>4</sub> ]]	1487.495	-0.8972	2.6966	0.48
[C <sub>2</sub> eim][Tf <sub>2</sub> N]	1845.795	-1.4624	6.036	0.25
[C <sub>10</sub> mim][Tf <sub>2</sub> N]	1546.69	-0.8939	-0.2304	0.77

Table 4. Coefficients of Equations 1 and 3, along with Standard Deviation of the Fit ( $\sigma$ ) and Average Absolute Relative Deviation (AARD) of Equation 1

ionic liquid	С	$b_1$ /MPa	$b_2{\boldsymbol{\cdot}}10^{-4}/\mathrm{MPa}{\boldsymbol{\cdot}}\mathrm{K}$	$\sigma/{\rm kg} \cdot {\rm m}^{-3}$	AARD/%
[C <sub>2</sub> mim][MeSO <sub>4</sub> ]	0.05567	-10.999	5.2634	0.42	0.027
[C <sub>2</sub> mim][EtSO <sub>4</sub> ]	0.08317	9.826	6.6572	0.03	0.022
[C <sub>2</sub> eim][Tf <sub>2</sub> N]	0.06313	-37.846	4.3762	0.19	0.011
[C <sub>10</sub> mim][Tf <sub>2</sub> N]	0.05446	-32.401	3.4109	1.34	0.057

Table 5. Calculated Densities by using Equation 11, Molar Mass (*M*), Molecular Volume (*V*), Experimental Average Density ( $\rho_{exptl-av}$ ), Calculated Average Density ( $\rho_{calcd-av}$ ), and Average Absolute Relative Deviation (AARD) of Ionic Liquids

	M	V/.	Å <sup>3</sup>	data	$\rho_{\mathrm{exptl-av}}$	$\rho_{\rm calcd-av}$	AARD
ionic liquid	$\overline{g \cdot mol^{-1}}$	cation	anion	points	$kg \cdot m^{-3}$	$kg \cdot m^{-3}$	%
[C <sub>2</sub> mim][MeSO <sub>4</sub> ]	222.26	182	106 <sup>a</sup>	80	1268.7	1266.0	0.38
[C2mim][EtSO4]	236.29	182	134 <sup>a</sup>	80	1226.2	1226.7	0.30
[C2eim][Tf2N]	405.34	210	248	80	1440.0	1451.9	0.83
[C <sub>10</sub> mim][Tf <sub>2</sub> N]	503.52	406	248	80	1259.6	1263.0	0.28

<sup>a</sup> Calculated in this work.

show that the Tait equation correlates well the pure ILs studied with an average absolute relative deviation (AARD) inferior to  $0.06 \ \%$ .

The following form of the Tait equation<sup>23</sup>

$$\rho = \frac{\rho(T, P = 0.1 \text{ MPa})}{\left\{1 - C \ln \frac{(B+P)}{(B+0.1)}\right\}}$$
(1)

where

$$\rho(T, P = 0.1 \text{ MPa}) = a_1 + a_2 T + a_3 T^2$$
(2)

was fitted to the density data. In eq 2,  $a_1$ ,  $a_2$ , and  $a_3$  were found by fitting to the experimental  $\rho(T, P = 0.1 \text{ MPa})$  and are given in Table 3.

Coefficient B is defined as

$$B = b_1 + \frac{b_2}{T} \tag{3}$$

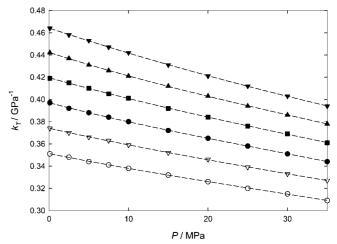
The coefficients *C*,  $b_1$ , and  $b_2$  were obtained by fitting the Tait equation to experimental data, and their values along with standard deviation  $\sigma$  of fit are given in Table 4. The standard deviation is defined by

$$\sigma = \left[\sum_{i=1}^{N_{\rm p}} \left(\rho_{\rm calcd} - \rho_{\rm exptl}\right)_i^2 / (N_{\rm p} - k)\right]^{1/2}$$
(4)

where  $N_p$  represents the number of data points ( $N_p = 80$ ) and k is the number of adjusted parameters (k = 3). The average absolute relative deviation (AARD) was defined as

$$AARD = \frac{\sum_{i=1}^{N_{p}} |(\rho_{calcd} - \rho_{exptl}) / \rho_{exptl}|_{i}}{N_{p}}$$
(5)

and is listed in Table 4. Figure 3 shows the good agreement between the experimental density data for  $[C_{2}eim][Tf_{2}N]$  and isotherms obtained with the Tait equation.



**Figure 4.** Isotherms for the isothermal compressibility of  $[C_2 \text{mim}][\text{EtSO}_4]$ :  $\bigcirc$ , 293.15 K;  $\Box$ , 298.15 K;  $\Delta$ , 303.15 K;  $\nabla$ , 313.15 K;  $\blacksquare$ , 333.15 K;  $\blacksquare$ , 353.15 K;  $\blacktriangle$ , 373.15 K;  $\blacktriangledown$ , 393.15 K. The curves were fitted to data.

The Tait equation is an integrated form of an empirical equation representative of the isothermal compressibility behavior versus pressure. The effect of pressure in density can be best described by the isothermal compressibility,  $\kappa_T$ , which is calculated according to the following expression

$$\kappa_T = -\frac{1}{V_{\rm m}} \left( \frac{\partial V_m}{\partial P} \right)_T = \frac{1}{\rho} \left( \frac{\partial \rho}{\partial P} \right)_T = \left( \frac{\partial \ln \rho}{\partial P} \right)_T \tag{6}$$

where  $\rho$  is the density and *P* the pressure at constant temperature, *T*. The isothermal compressibilities can be calculated using eqs 1 and 6

$$\kappa_T = \left(\frac{C}{B+P}\right) \left(\frac{\rho}{\rho(T,P=0.1 \text{ MPa})}\right) \tag{7}$$

For illustration purposes, the isothermal compressibilities of  $[C_2mim][EtSO_4]$  are shown in Figure 4. The ILs become more compressible with increasing temperatures and less compressible with increasing pressure. The calculated values of  $\kappa_T$  are presented in the Supporting Information, as Table ST1. In the studied range of temperature, (293.15 to 393.15) K, and pressure, (0.10 to 35.00) MPa, the isothermal compressibility ranges in GPa<sup>-1</sup> are 0.28 to 0.45, 0.31 to 0.46, 0.44 to 0.86, and 0.47 to 1.00, respectively, for  $[C_2mim][MeSO_4]$ ,  $[C_2mim][EtSO_4]$ ,  $[C_2eim][Tf_2N]$ , and  $[C_{10}mim][Tf_2N]$ .

The isobaric expansivity,  $\alpha_P$ , is defined as

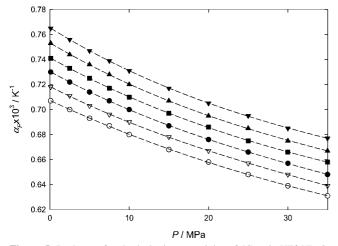
$$\alpha_{P} = \frac{1}{V_{m}} \left( \frac{\partial V_{m}}{\partial T} \right)_{P} = -\frac{1}{\rho} \left( \frac{\partial \rho}{\partial T} \right)_{P} = -\left( \frac{\partial \ln \rho}{\partial T} \right)_{P}$$
(8)

and the following expression is derived from the Tait equation, eq 1

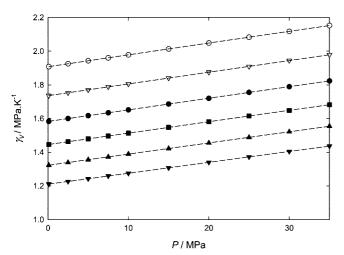
$$\alpha_{P} = -\left\{ \frac{\left[\frac{d\rho(T, P = 0.1 \text{ MPa})/dT\right]}{\rho(T, P = 0.1 \text{ MPa})} \right\} + C\left\{ \frac{\frac{dB}{dT}(P - 0.1)}{\left[1 - C \ln\left(\frac{B+P}{B+0.1}\right)\right](B + 0.1)(B+P)} \right\}$$
(9)

where  $dB/dT = -b_2/T^2$ . The values of  $\alpha_P$  are presented in the Supporting Information, as Table ST2, for all investigated ILs.

The  $\alpha_p$  values for  $[C_{10}mim][Tf_2N]$  are shown in Figure 5. Studied ILs do not expand notably with temperature. However, there are some ILs that present a more significant



**Figure 5.** Isotherms for the isobaric expansivity of  $[C_{10}mim][Tf_2N]: \bigcirc$ , 293.15 K;  $\Box$ , 298.15 K;  $\Delta$ , 303.15 K;  $\bigtriangledown$ , 313.15 K;  $\bullet$ , 333.15 K;  $\blacksquare$ , 353.15 K;  $\bigstar$ , 373.15 K;  $\blacktriangledown$ , 393.15 K. The curves were fitted to data.



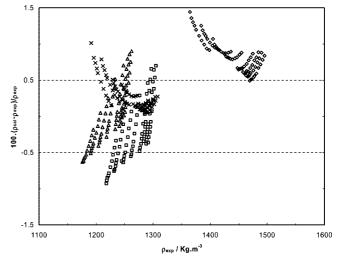
**Figure 6.** Isotherms for the thermal pressure coefficient of  $[C_2mim][MeSO_4]$ : O, 293.15 K;  $\Box$ , 298.15 K;  $\Delta$ , 303.15 K;  $\nabla$ , 313.15 K;  $\blacksquare$ , 333.15 K;  $\blacksquare$ , 353.15 K;  $\blacktriangle$ , 373.15 K;  $\blacktriangledown$ , 393.15 K. The curves were fitted to data.

decreases in  $\alpha_p$  with temperature,<sup>36</sup> but in fact, ILs do not expand markedly with temperature and have  $\alpha_p$  values lower than most organic liquids and similar to that of water. From eq 9, the isobaric expansivity, for the studied range of temperature (293.15 to 393.15) K and pressure (0.10 to 35.00) MPa, is {(0.5 to 0.6)  $\cdot 10^{-3}$ , (0.5 to 0.6)  $\cdot 10^{-3}$ , (0.6 to 0.8)  $\cdot 10^{-3}$ , and (0.6 to 0.8)  $\cdot 10^{-3}$ } K<sup>-1</sup>, respectively, for [C<sub>2</sub>mim][MeSO<sub>4</sub>], [C<sub>2</sub>mim][EtSO<sub>4</sub>], [C<sub>2</sub>eim][Tf<sub>2</sub>N], and [C<sub>10</sub>mim][Tf<sub>2</sub>N].

The thermal pressure coefficient,  $\gamma_V$ , may be calculated according to

$$\gamma_V = \frac{\alpha_p}{\kappa_T} \tag{10}$$

The thermal pressure coefficients as a function of pressure obtained for investigated ILs are shown in Figure 6 for  $[C_2mim][MeSO_4]$ . The  $\gamma_V$  decreases with temperature and increases slightly with pressure for all the ILs studied (see Supporting Information, Table ST3). We have applied the propagation law of errors at eqs 7, 9, and 10 taking into account the uncertainties in the parameters *c*, *b*<sub>1</sub>, and *b*<sub>2</sub> of the Tait equation, eq 1, and *a*<sub>1</sub> to *a*<sub>3</sub> of eq 2, and also in the temperature and pressure. We have concluded that the



**Figure 7.** Relative deviations between the experimental and calculated density using eq 11 as a function of experimental density of ionic liquids:  $\Box$ , [C<sub>2</sub>mim][MeSO<sub>4</sub>];  $\Delta$ , [C<sub>2</sub>mim][EtSO<sub>4</sub>];  $\diamond$ , [C<sub>2</sub>eim][Tf<sub>2</sub>N]; ×, [C<sub>10</sub>mim]-[Tf<sub>2</sub>N].

uncertainty in  $\kappa_T$  was of the order of  $\pm 0.05$  GPa<sup>-1</sup>, while for  $\alpha_p$  it will be of the order of  $\pm 5 \cdot 10^{-4}$  K<sup>-1</sup>. The uncertainty in  $\gamma_V$  was determined with the law of propagation of errors from the uncertainty in  $\alpha_p$  and  $\kappa_T$ . We have obtained a maximum uncertainty of  $\pm 1$  MPa·K<sup>-1</sup>, which is almost of the same magnitude of the thermal pressure coefficient.

Correlation by the Gardas and Coutinho Estimation Method. Recently, we proposed a new method<sup>25</sup> for the estimation of ionic liquid densities in a wide range of temperatures, (273.15 to 393.15) K, and pressures, (0.10 to 100) MPa, according to eq 11

$$\rho = \frac{M}{NV(a+bT+cP)} \tag{11}$$

where  $\rho$  is the density in kg·m<sup>-3</sup>; *M* is molar mass in kg·mol<sup>-1</sup>; N is the Avogadro constant; V is the molecular volume in  $Å^3$ ; T is the temperature in K; and P is the pressure in MPa. The coefficients a, b, and c were estimated by fitting eq 11 to our previously published experimental data.8,10 A total amount of ca. 800 density data points were used. The values of coefficient *a*, *b*, and *c* obtained are  $8.005 \cdot 10^{-1} \pm 2.333 \cdot 10^{-4}$ , (6.652  $\cdot 10^{-4}$ )  $\pm 6.907 \cdot 10^{-7}$ ) K<sup>-1</sup>, and (-5.919  $\cdot 10^{-4} \pm 2.410 \cdot 10^{-6}$ ) MPa<sup>-1</sup>, respectively, at 95 % confidence level. The average absolute relative deviation of calculated densities from the experimental densities is 0.29 %. The experimental density (Y) for the ILs used to obtain coefficients a, b, and c of eq 11 is essentially identical to its calculated density (X):  $Y = (0.9998 \pm 0.0003)X$ (correlation coefficient:  $R^2 = 0.9989$ , at 95 % level of confidence). For the density calculation, volume parameters (V)of ions and groups were either directly taken from the literature<sup>37</sup> or calculated following the Jenkins procedure.<sup>38</sup> The volume of the sulfate anion  $(SO_4^{-2})$  is estimated in this work as 71 Å<sup>3</sup> by minimizing the objective function (O.F.)

$$O.F. = \frac{\sum_{i=1}^{N_{p}} |(\rho_{calcd} - \rho_{exptl}) / \rho_{exptl}|_{i}}{N_{p}}$$
(12)

where  $N_p$  represents the number of data points ( $N_p = 160$ ) of  $[C_2 \text{mim}][\text{MeSO}_4]$  and  $[C_2 \text{mim}][\text{EtSO}_4]$ .

Table 6. Sanchez–Lacombe Equation of State Parameters, Characteristic Pressure ( $P^*$ ), Temperature ( $T^*$ ), and Density ( $\rho^*$ ), Segment Number of One Molecule (r), Segment Interaction Energy ( $\varepsilon^*$ ), Segment Volume ( $v^*$ ), and Hard Core Volume ( $v^* \times r$ ) and Average Absolute Relative Deviation (AARD) of Ionic Liquids

	<i>P</i> *	<i>T</i> *	$\rho^*$		£*	ν*	$\nu^* \times r$	AARD
ionic liquid	MPa	Κ	$kg \cdot m^{-3}$	r	$J \cdot mol^{-1}$	$cm^3 \cdot mol^{-1}$	$cm^3 \cdot mol^{-1}$	%
[C <sub>2</sub> mim][MeSO <sub>4</sub> ]	575.2	685.7	1359.2	16.5	5701	9.91	163.5	0.13
$[C_2 mim][EtSO_4]$	574.3	675.2	1316.9	18.4	5614	9.78	179.4	0.15
$[C_2 eim][Tf_2N]$	498.0	603.2	1580.4	25.5	5015	10.07	256.5	0.18
[C <sub>10</sub> mim][Tf <sub>2</sub> N]	544.7	598.4	1382.0	39.9	4975	9.13	364.3	0.20

The densities of  $[C_2mim][MeSO_4]$ ,  $[C_2mim][EtSO_4]$ ,  $[C_2eim][Tf_2N]$ , and  $[C_{10}mim][Tf_2N]$  were calculated by using eq 11. Calculated densities with the average absolute relative deviation (AARD) of studied ILs are shown in Table 5. As shown in Figure 7, the experimental densities (*Y*) of  $[C_2mim][MeSO_4]$ ,  $[C_2mim][EtSO_4]$ ,  $[C_2eim][Tf_2N]$ , and  $[C_{10}mim][Tf_2N]$  display a very good agreement with the predicted densities (*X*) from eq (11):  $Y = (0.9972 \pm 0.0003)X$ (correlation coefficient:  $R^2 = 0.9947$ , at a 95 % level of confidence).

These results (Table 5 and Figure 7) show that the density estimation method proposed by Gardas and Coutinho<sup>25</sup> can be applied with confidence to ionic liquids from families other than those used on the development of the correlation.

Correlation by the Sanchez–Lacombe Equation of State. The Sanchez–Lacombe equation of state (S–L EoS) is based on lattice fluid theory<sup>24</sup> and can be used to predict phase equilibria from the pure component parameters. This equation of state is given by the following equations

$$\tilde{\rho}^{2} + \tilde{P} + \tilde{T} \left[ \ln(1 - \tilde{\rho}) + \left(1 - \frac{1}{r}\right) \tilde{\rho} \right] = 0$$
  

$$\tilde{\rho} \equiv \rho / \rho^{*} \quad \tilde{P} \equiv P / P^{*} \quad \tilde{T} \equiv T / T^{*}$$
  

$$\varepsilon^{*} = RT^{*} \quad \upsilon^{*} = RT^{*} / P^{*} \quad r = MP^{*} / RT^{*} \rho^{*}$$
(13)

where  $P^*$ ,  $T^*$ , and  $\rho^*$  represent characteristic parameters of the EOS; *M* is molar mass; and *R* is the gas constant. For studied ILs, the characteristic parameters were determined by fitting the equation of state (eq 13) to the experimental pressure-volume-temperature (PVT) data in the region from (293 to 393) K and from (0.1 to 35) MPa.

The S-L EoS parameters and the segment number of one molecule (*r*), segment interaction energy ( $\varepsilon^*$ ), segment volume ( $\nu^*$ ), and hard core volume ( $\nu^* \times r$ ) calculated from the S-L EoS parameters are summarized in Table 6. Figure 8 shows the relative deviations between the experimental and calculated density data as a function of experimental densities, using eq 13, for studied ILs. The maximum deviation between calculated and experimental density data was within 0.53 % for all ionic liquids.

The free volume  $(V_f)$  can be calculated as

$$V_{\rm f} = V_{\rm m} - (v^* \times r) \tag{14}$$

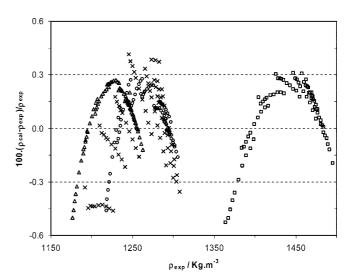
where  $V_{\rm m}$  is the molar volume of IL at given *T* and *P*. This equation was applied by Machida et al.<sup>7</sup> to determine the free volume of ILs. Figure 9 shows free volumes, calculated using eq 14, of studied ILs as a function of temperature at (0.1 and 35) MPa. As expected, the free volume of studied ILs increased with temperature and decreased with pressure. For the studied ILs, free volume is increasing showing the trend:  $[C_2mim][MeSO_4] < [C_2mim][EtSO_4] < [C_2eim][Tf_2N] < [C_{10}mim][Tf_2N].$ 

# Conclusions

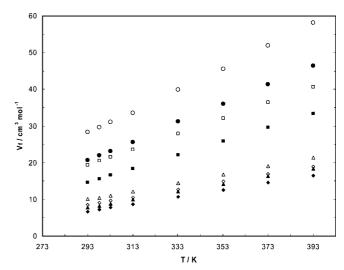
Experimental density data for four pure ILs in the temperature range (293.15 to 393.15) K and pressure range

(0.10 to 35.00) MPa are presented. Density results can be explained by structural variations in the cation and anion. From the experimental data, a proportional molar volume increase with the alkyl chain length of the cation of ILs was observed. The liquid densities were correlated with the Tait equation<sup>23</sup> that has shown to describe well all the pure ILs studied with an average absolute relative deviation (AARD) less than 0.06 %.

The experimental results were also used to estimate some derivative thermodynamic properties such as the isothermal compressibility, the isobaric expansivity, and the thermal



**Figure 8.** Relative deviations between the experimental and calculated density using eq 13 as a function of experimental density of ionic liquids:  $\Box$ , [C<sub>2</sub>mim][MeSO<sub>4</sub>];  $\Delta$ , [C<sub>2</sub>mim][EtSO<sub>4</sub>];  $\bigcirc$ , [C<sub>2</sub>eim][Tf<sub>2</sub>N];  $\times$ , [C<sub>10</sub>mim]-[Tf<sub>2</sub>N].



**Figure 9.** Free volume of investigated ionic liquids,  $[C_2mim][MeSO_4]: \diamond$ , 0.1 MPa;  $\blacklozenge$ , 35 MPa.  $[C_2mim][EtSO_4]: \Delta$ , 0.1 MPa;  $\blacklozenge$ , 35 MPa.  $[C_2eim][Tf_2N]: \Box$ , 0.1 MPa;  $\blacksquare$ , 35 MPa.  $[C_{10}mim][Tf_2N]: \bigcirc$ , 0.1 MPa;  $\blacklozenge$ , 35 MPa.

pressure coefficient of the studied ILs that are difficult to obtain by direct measurements at extreme conditions of pressure and temperature.

The Gardas and Coutinho<sup>25</sup> method for the prediction of ionic liquid densities was tested against the measured densities and results showing that it applies to other ionic liquid families than those used for its development.

The Sanchez–Lacombe equation of state was used to correlate the density data of studied ionic liquids. The maximum deviation between calculated and experimental density data was within 0.53 % for all ionic liquids.

#### **Supporting Information Available:**

An additional three tables. This material is available free of charge via the Internet at http://pubs.acs.org.

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